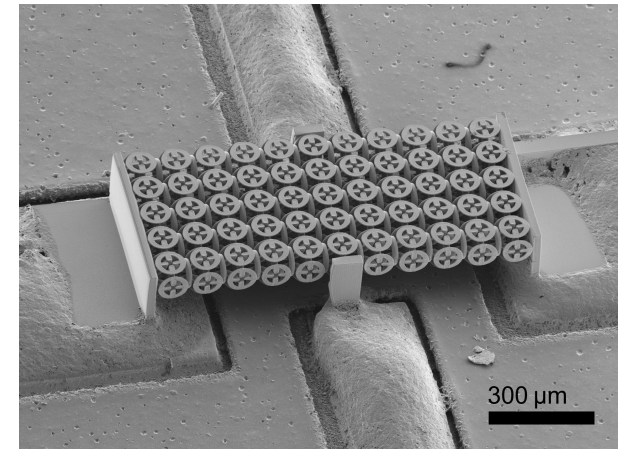
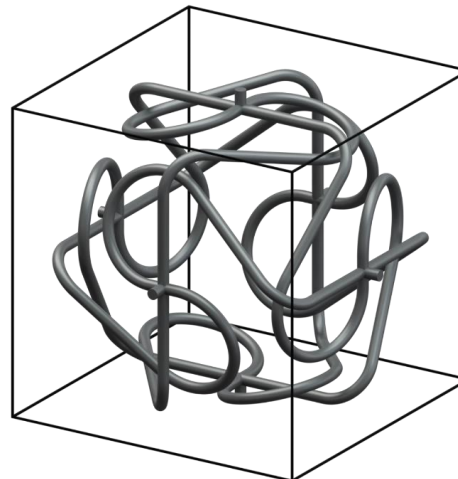
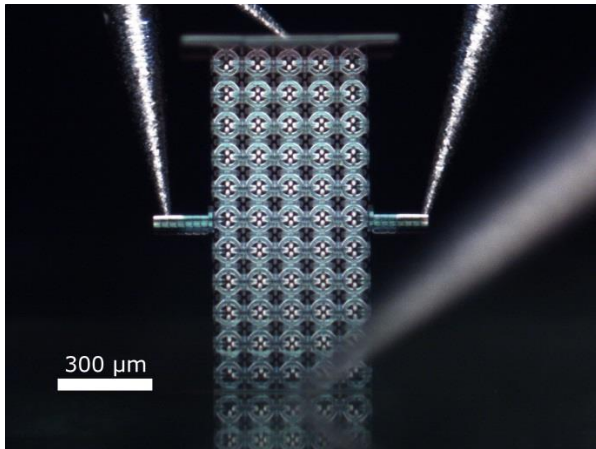


# The Hall effect in composites

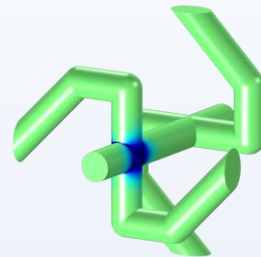
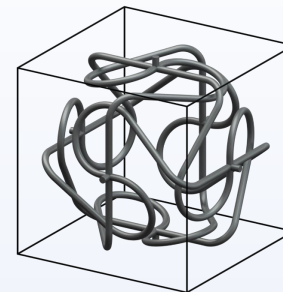
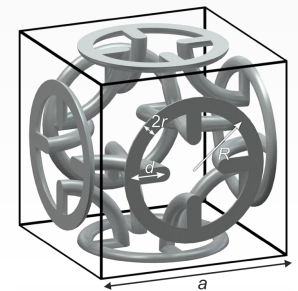
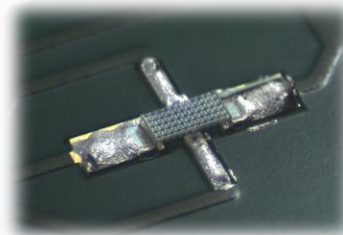
**Christian Kern**, KIT, Germany  
Graeme W Milton, University of Utah, USA  
Muamer Kadic, FEMTO-ST, France  
Martin Wegener, KIT, Germany  
christian.kern@kit.edu

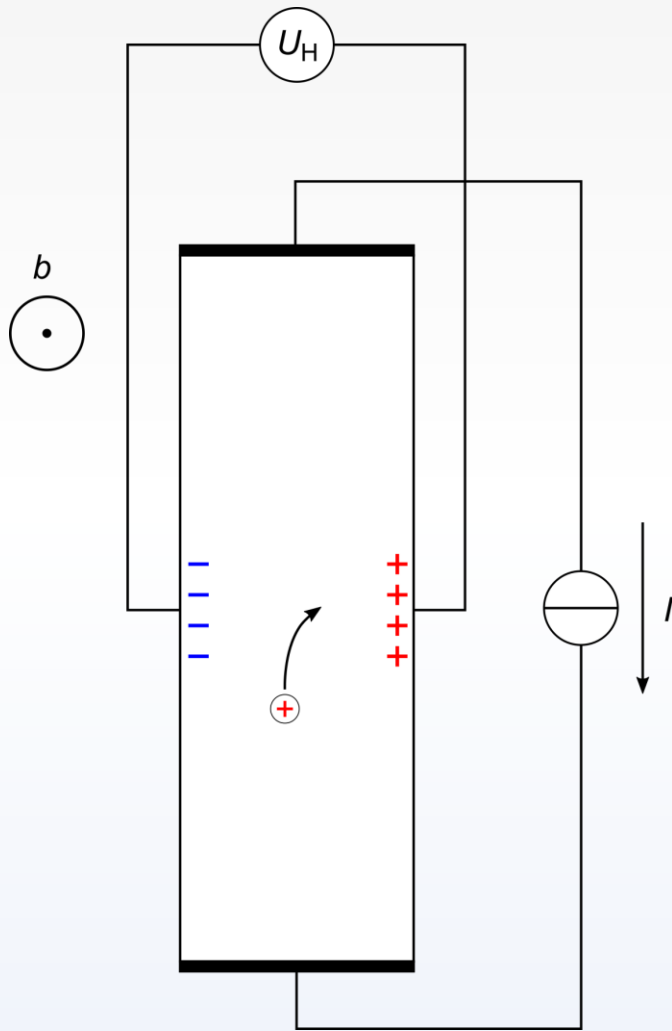


# Outline

- Theory and homogenization of the Hall effect
- Sign-inversion of the effective Hall coefficient in chainmail-inspired composites
- Symmetry considerations
- Current state of experiments
- Novel structures showing a sign-inversion of the effective Hall coefficient
- Anisotropic structures

$$\text{Cof}(\sigma_0^*) \mathbf{A}_H^* = \langle \text{Cof}(\sigma_0 \nabla \Phi)^T \mathbf{A}_H \rangle$$





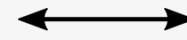
$$U_H = \frac{1}{nq} \frac{Ib}{h} = A_H \frac{Ib}{h}$$

negative charge carriers (e.g. n-type semiconductor)

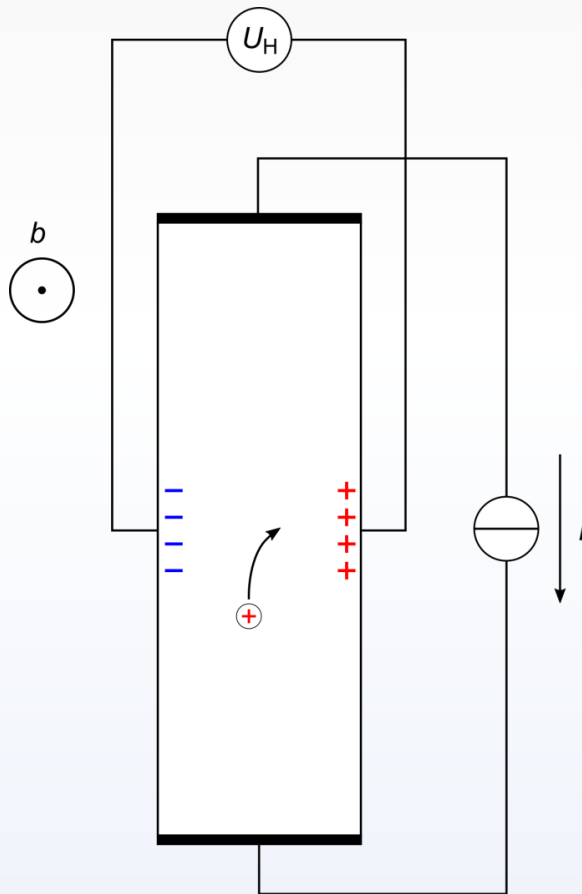


negative  $A_H$

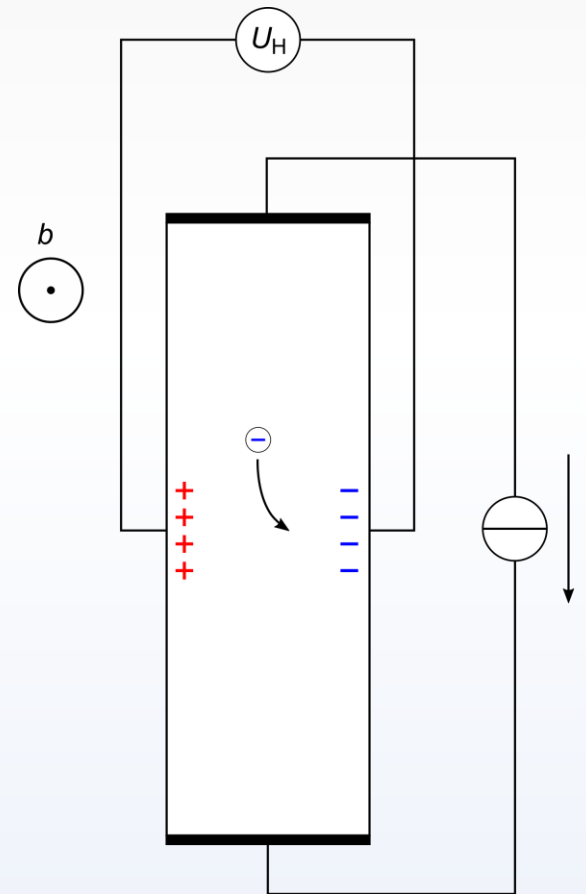
positive charge carriers (e.g. p-type semiconductor)



positive  $A_H$



positive Hall coefficient



negative Hall coefficient

Conductivity problem  $\nabla \cdot \mathbf{j} = 0$ ,  $\mathbf{j} = \boldsymbol{\sigma}(\mathbf{b}) \mathbf{e}$ ,  $\nabla \times \mathbf{e} = 0$

From Onsager's principle:  $\boldsymbol{\sigma}(\mathbf{b}) = \boldsymbol{\sigma}(-\mathbf{b})^\top$

### Small magnetic fields

Expansion in orders of the magnetic flux density:  $\boldsymbol{\sigma}(\mathbf{b}) = \boldsymbol{\sigma}_0 + \boldsymbol{\sigma}_1(\mathbf{b})$

$$\boldsymbol{\sigma}_0^\top = \boldsymbol{\sigma}_0$$

$$\boldsymbol{\sigma}_1(\mathbf{b})^\top = -\boldsymbol{\sigma}_1(\mathbf{b})$$

$$\begin{aligned} \boldsymbol{\sigma} &= \boldsymbol{\sigma}_0 + \mathcal{E}(\mathbf{S}\mathbf{b}) \\ \boldsymbol{\rho} &= \boldsymbol{\rho}_0 + \mathcal{E}(\mathbf{A}_H\mathbf{b}) \end{aligned} \quad \begin{array}{l} \curvearrowright \\ \mathbf{S} = -\text{Cof}(\boldsymbol{\sigma}_0) \mathbf{A}_H \end{array}$$

$$\text{Cof}(\mathbf{A})_{12} = (-1)^{1+2} \det \begin{pmatrix} \overline{A_{11}} & \overline{A_{12}} & \overline{A_{13}} \\ A_{21} & A_{22} & A_{23} \\ A_{31} & A_{32} & A_{33} \end{pmatrix} = -(A_{21}A_{33} - A_{31}A_{23})$$

L. D. Landau and E. M. Lifshitz, *Electrodynamics of Continuous Media*, 1960

M. Briane and G.W. Milton, *Arch. Ration. Mech. Anal.* **193**, 715 (2009)

## Isotropic materials

$$\boldsymbol{\sigma}_0 = \sigma_0 \mathbf{I} \quad \mathbf{A}_H = A_H \mathbf{I}$$

$$\mathbf{b}\text{-field along } \hat{\mathbf{z}}: \quad \boldsymbol{\sigma}(\mathbf{b}) = \begin{pmatrix} \sigma_0 & \sigma_0 A_H b_z & 0 \\ -\sigma_0 A_H b_z & \sigma_0 & 0 \\ 0 & 0 & \sigma_0 \end{pmatrix}$$

Effective Hall tensor

$$\mathbf{A}_H(\mathbf{x}) \longrightarrow \mathbf{A}_H^*$$

**Small magnetic fields**

$\mathbf{A}_H^*$  can be obtained from a solution of the zero magnetic-field problem

$$\text{Cof}(\boldsymbol{\sigma}_0^*) \mathbf{A}_H^* = \langle \text{Cof}(\boldsymbol{\sigma}_0 \nabla \Phi)^\top \mathbf{A}_H \rangle$$

$$\Phi = (\phi_1, \phi_2, \phi_3)^\top \quad \nabla \cdot (\boldsymbol{\sigma} \nabla \Phi) = 0$$

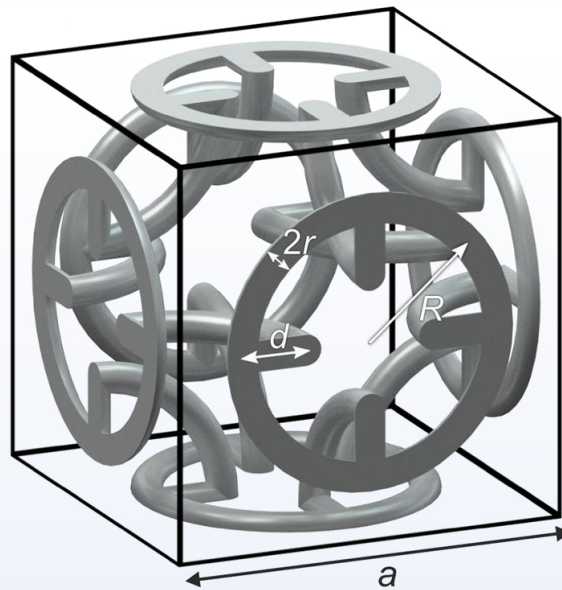
Isotropic case first considered by D. Bergman

$$\mathbf{A}_H^* = \left\langle \left( \tilde{J}_{11} \tilde{J}_{22} - \tilde{J}_{21} \tilde{J}_{12} \right) \mathbf{A}_H(\mathbf{x}) \right\rangle$$

M. Briane and G.W. Milton, Arch. Ration. Mech. Anal. **193**, 715 (2009)

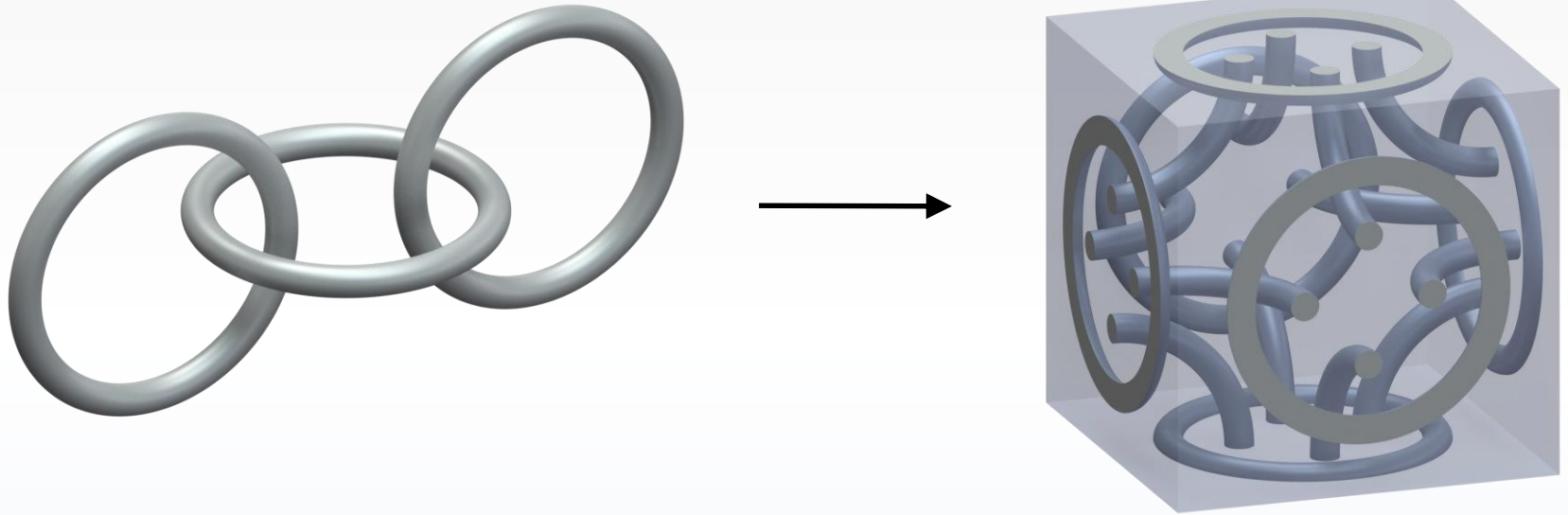
D. Bergman, in *Percolation Structures and Processes*, eds. G. Deutscher, R. Zallen, and J. Adler, 1983, pp. 297

# Sign-inversion of the effective Hall coefficient





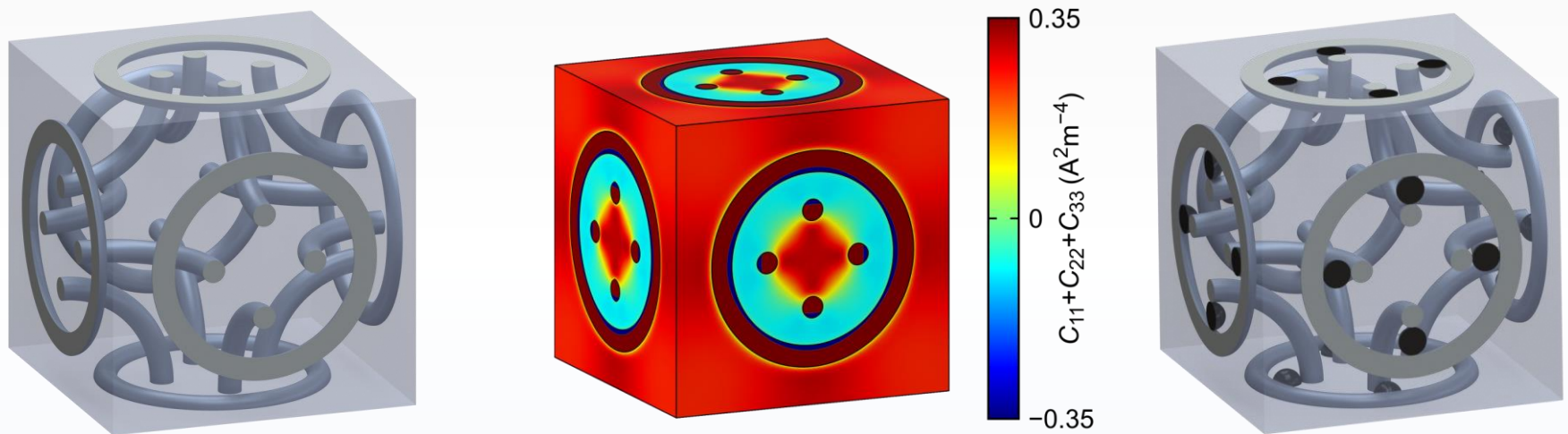
Change of sign of the determinant



$$\text{Cof}(\boldsymbol{\sigma}_0^*) \mathbf{A}_H^* = \langle \text{Cof}(\boldsymbol{\sigma}_0 \nabla \Phi)^\top \mathbf{A}_H \rangle$$

M. Briane, G. W. Milton, and V. Nesi, Arch. Ration. Mech. Anal. **173**, 133 (2004)

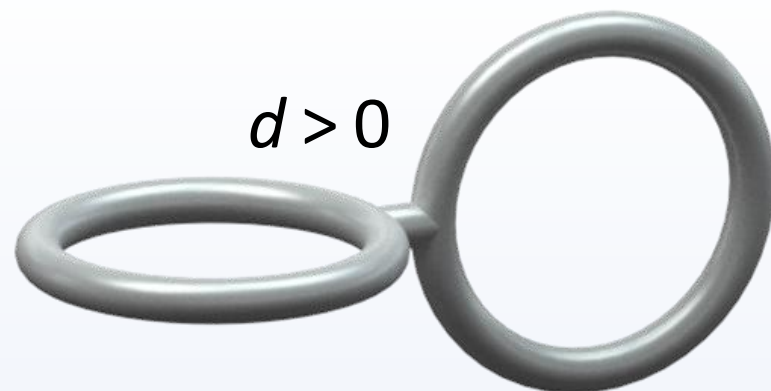
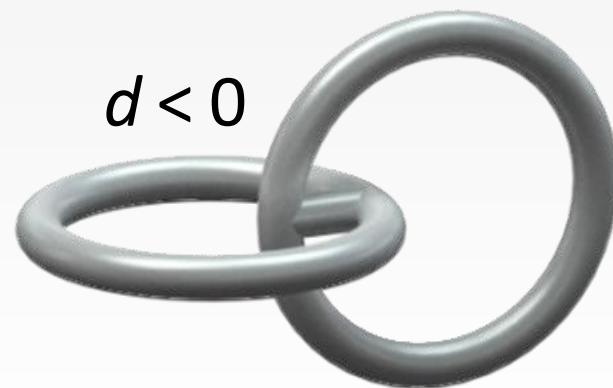
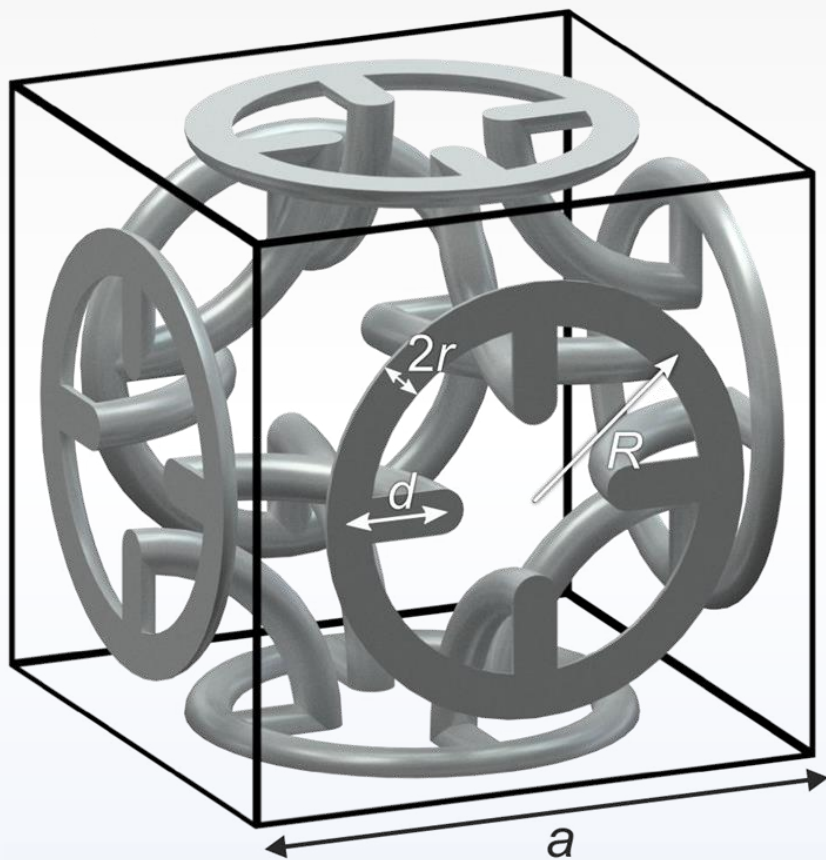
3D structure suggested by chainmail-artist Dylan Whyte

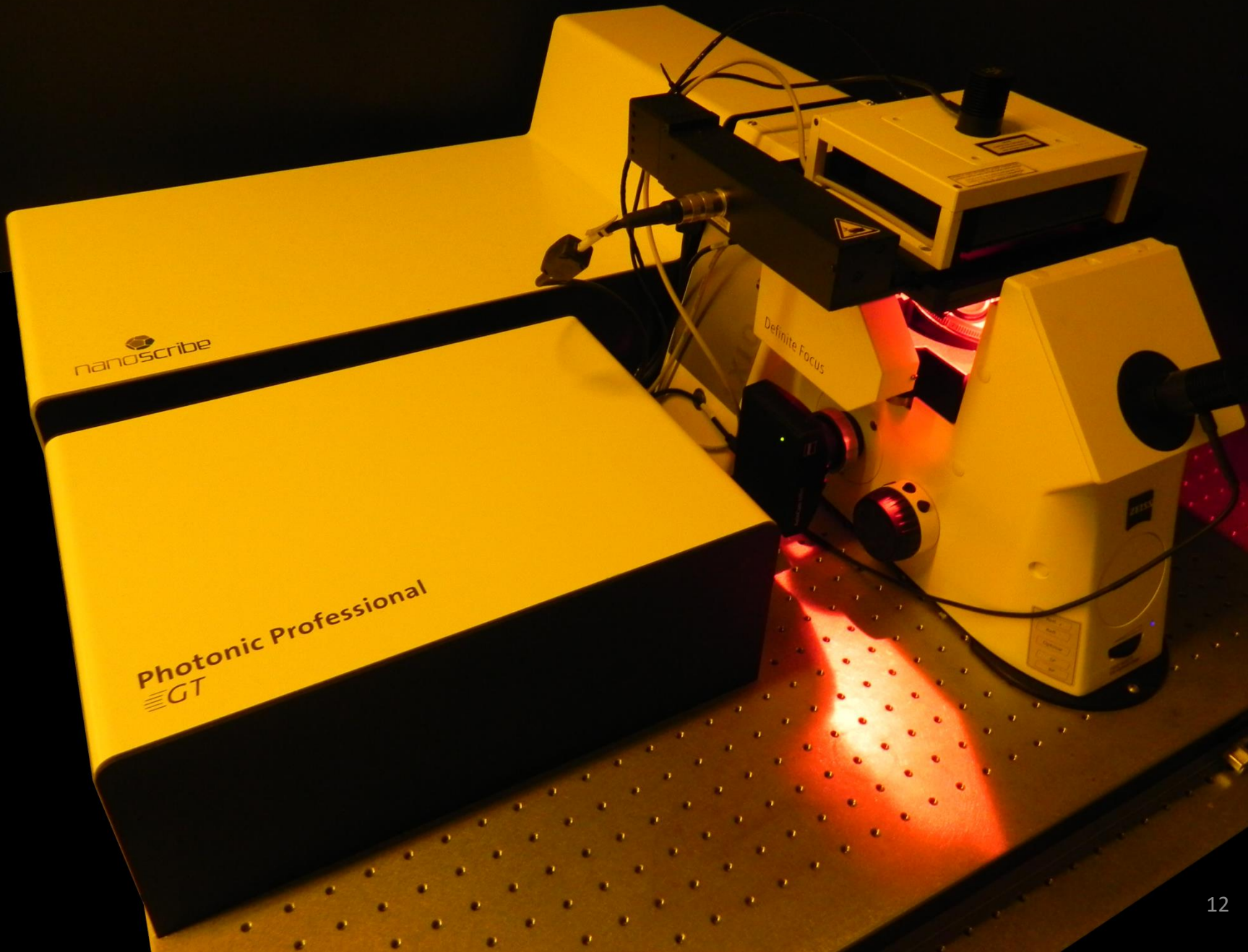


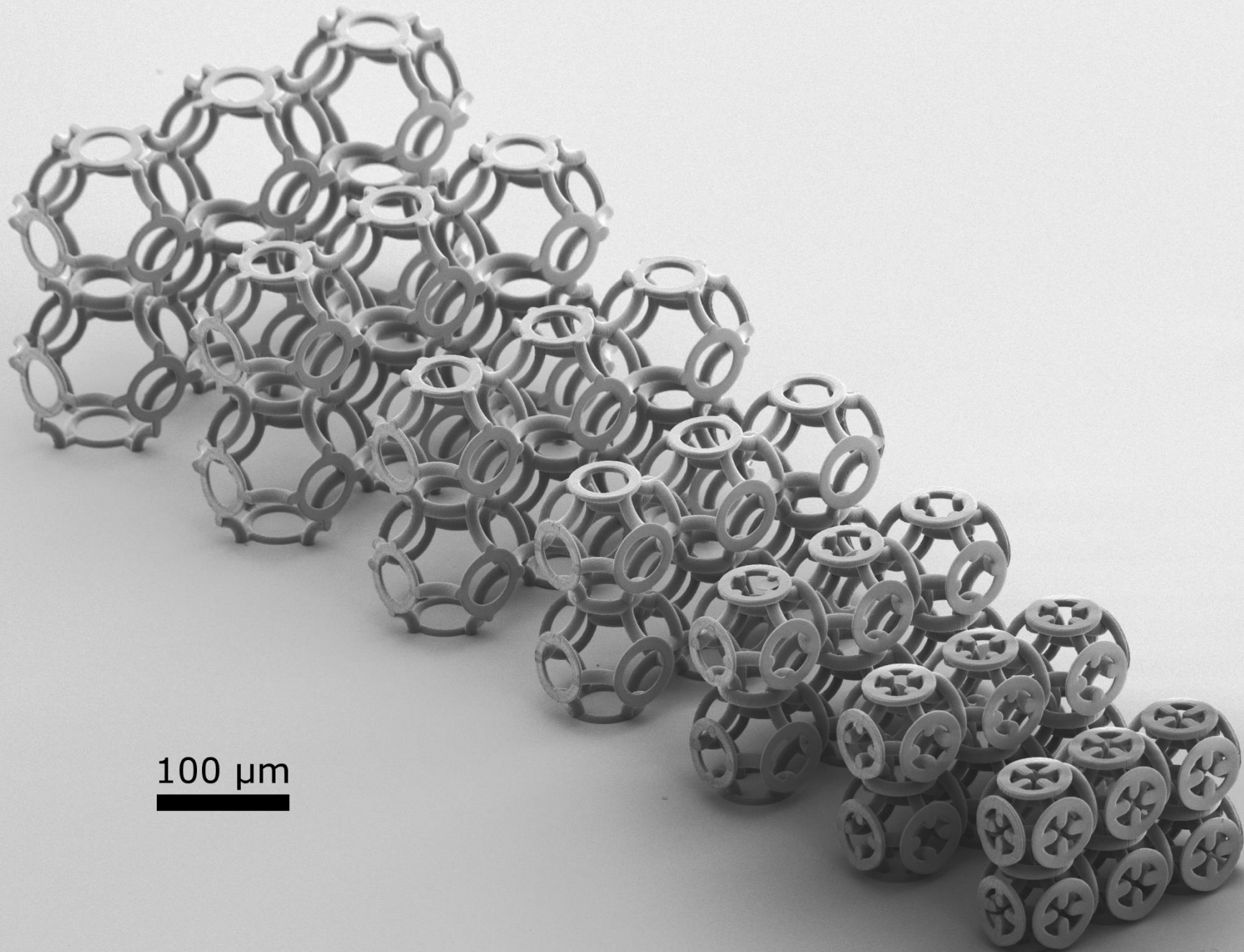
$$\text{Cof}(\boldsymbol{\sigma}_0^*) \mathbf{A}_H^* = \langle \text{Cof}(\boldsymbol{\sigma}_0 \nabla \Phi)^\top \mathbf{A}_H \rangle$$

M. Briane and G. W. Milton, Arch. Ration. Mech. Anal. **193**, 715 (2009)

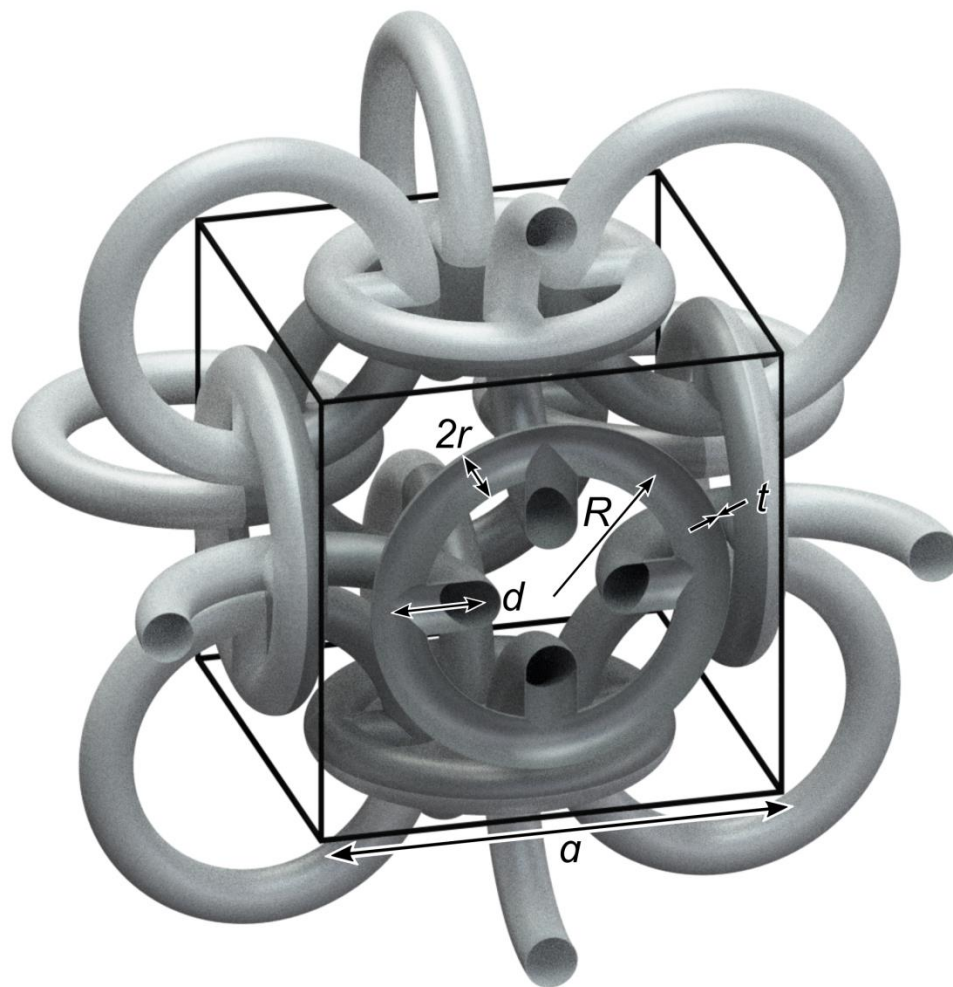
C. Kern *et al.*, arXiv:1806.04914 [cond-mat.mes-hall]

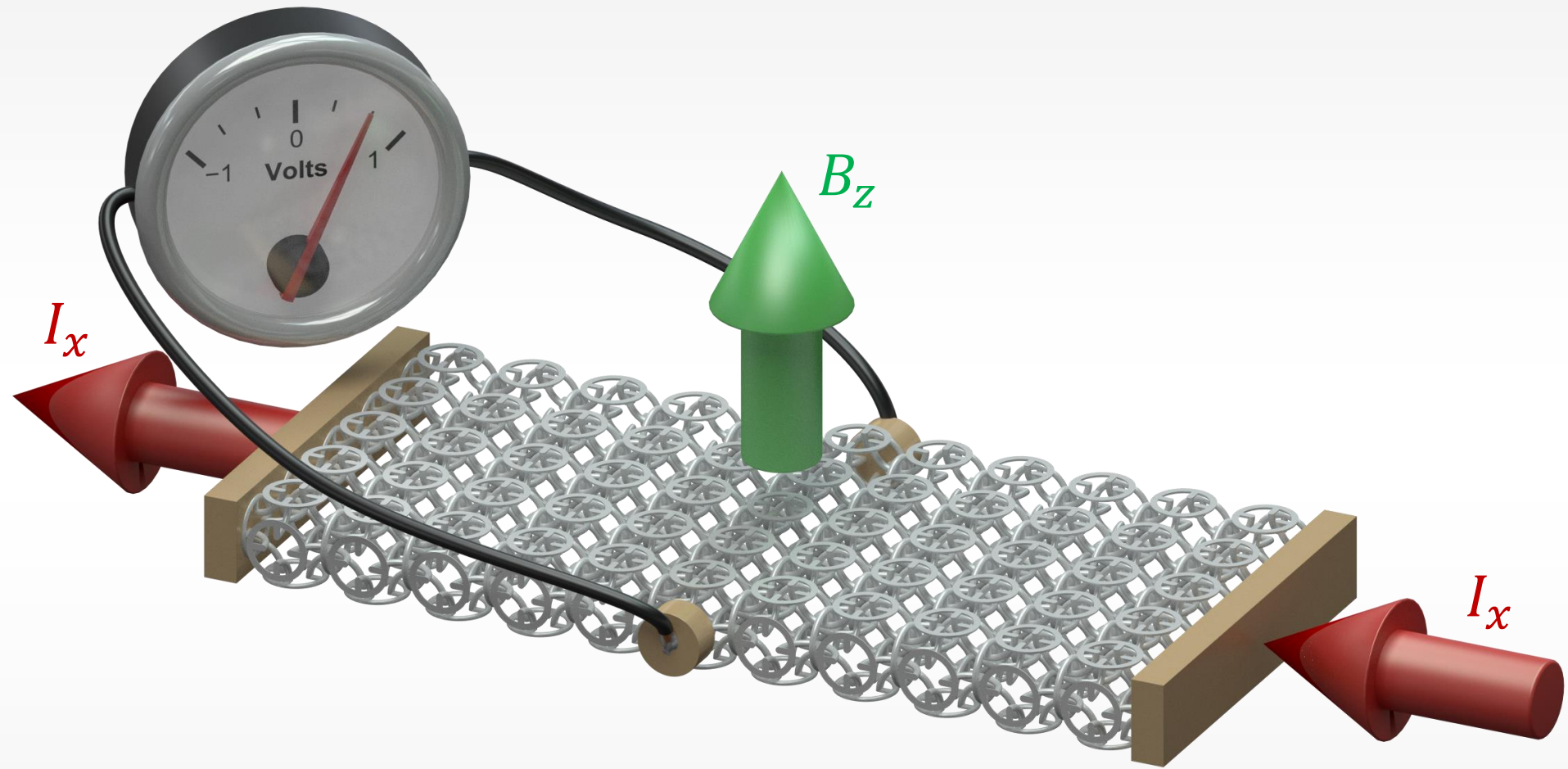


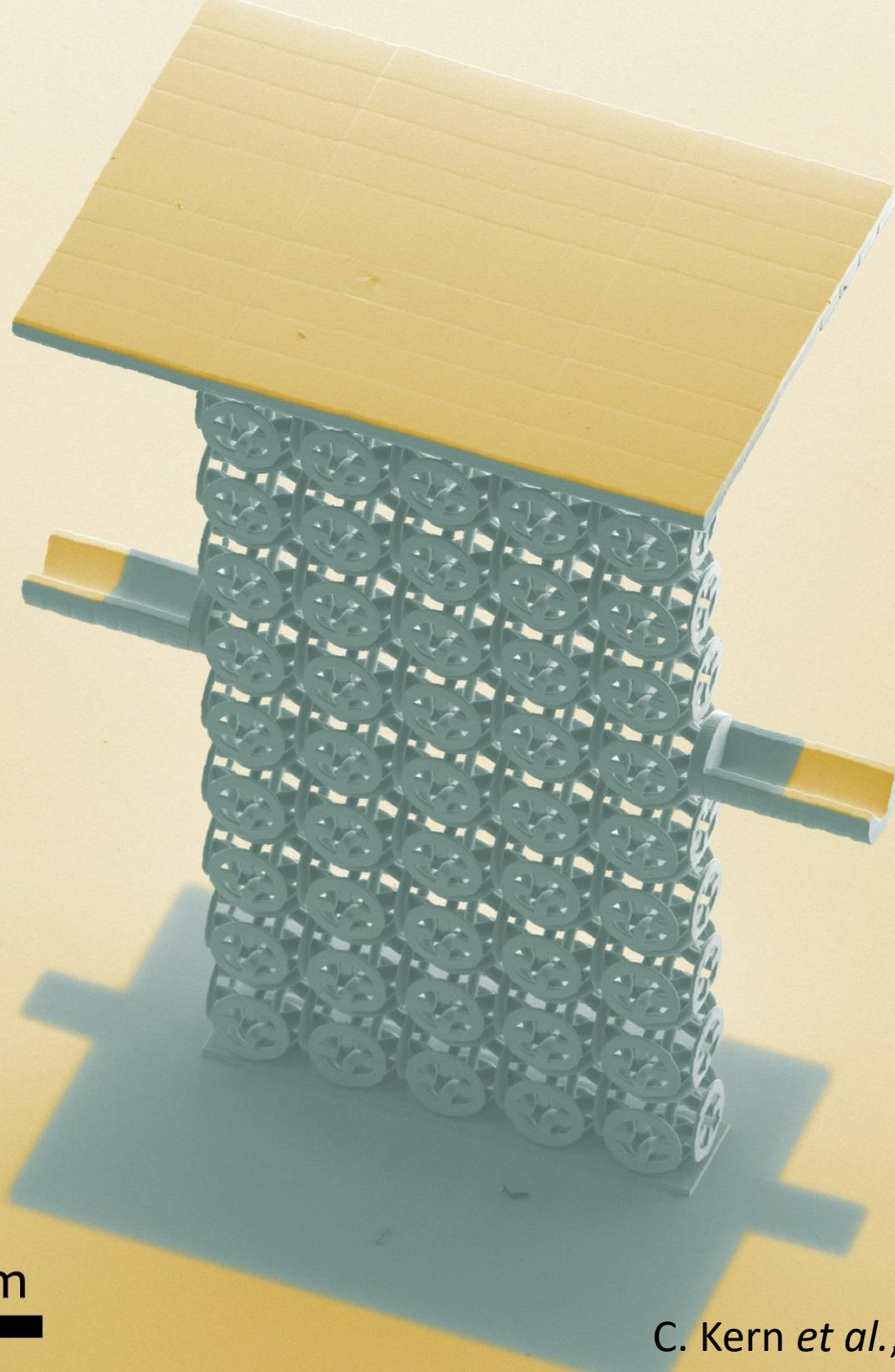




100  $\mu\text{m}$

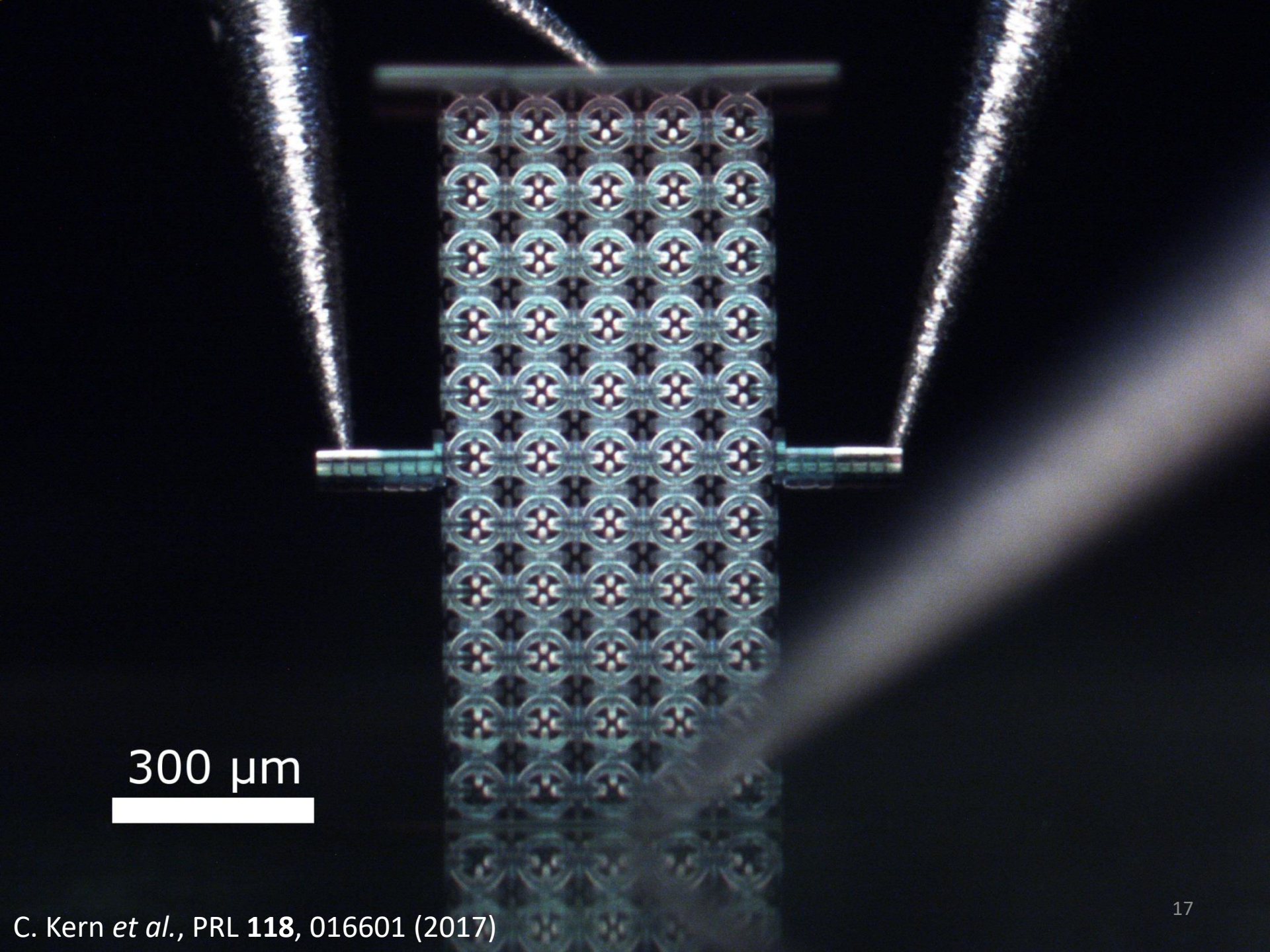




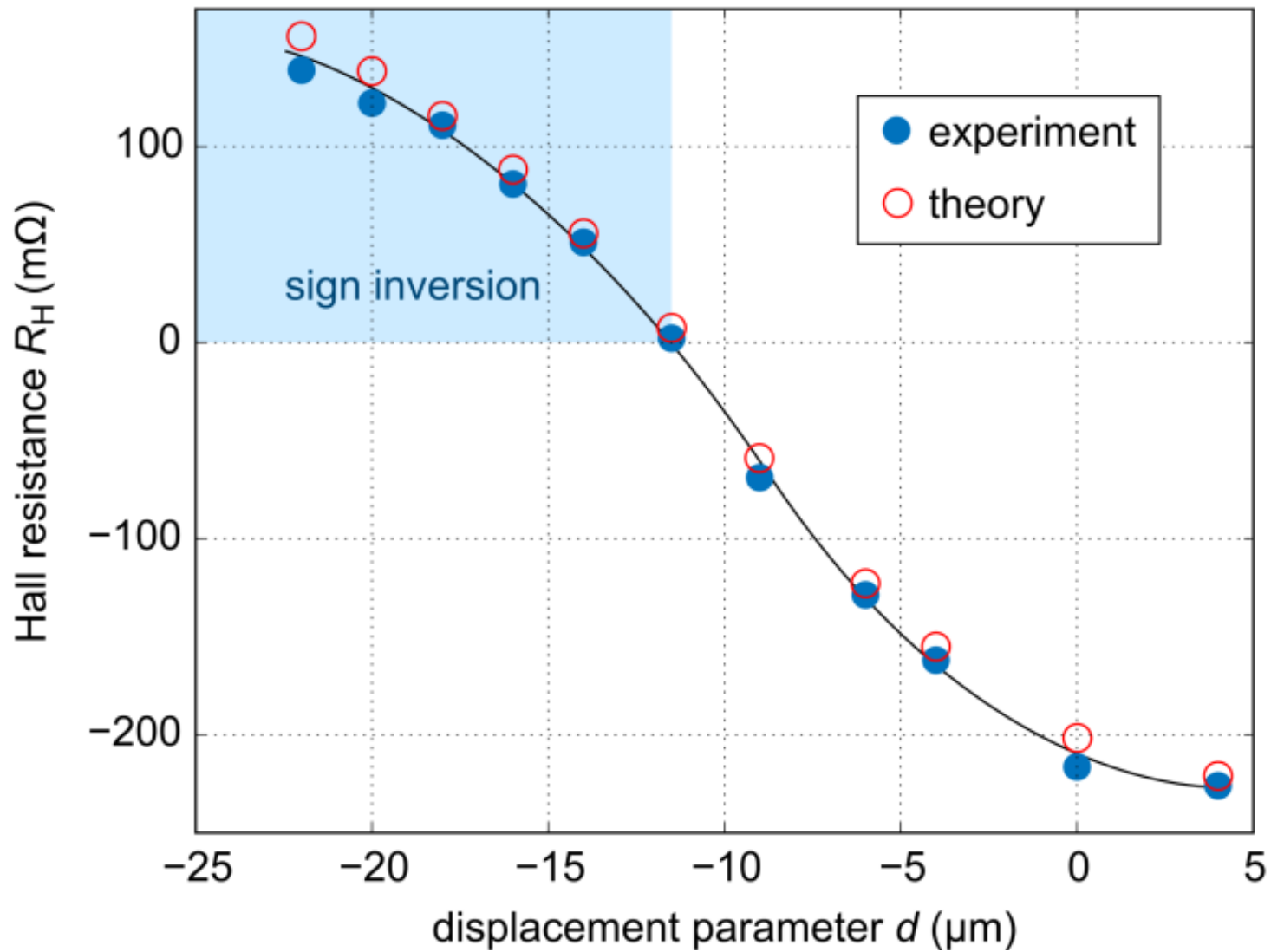


200  $\mu\text{m}$

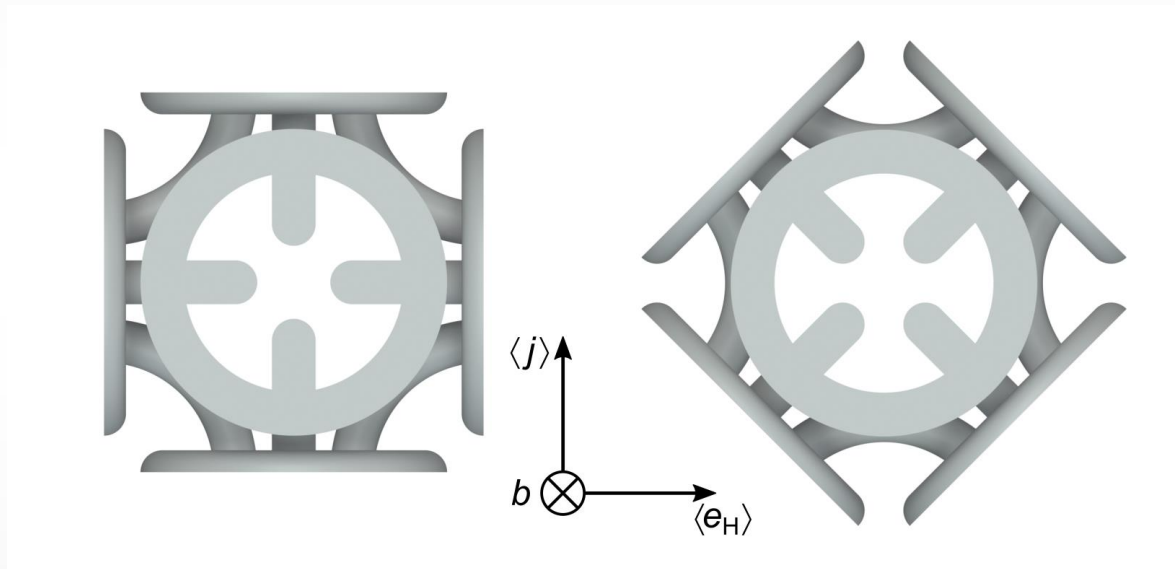


A scanning electron micrograph (SEM) of a photonic crystal slab. The slab is rectangular and features a periodic lattice of circular holes. A central waveguide is formed by a row of larger holes. The slab is mounted on a substrate, and a thin layer is visible on top. The image is illuminated from the left, creating a bright glow on the left side of the slab.

300  $\mu\text{m}$



What happens for other unit cell orientations?



# Symmetry considerations

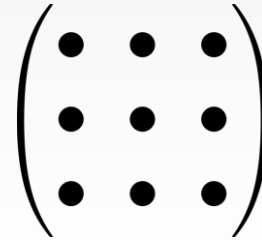
## Neumann's principle

*If a system has a certain group of symmetry operations then any physical observable of that system must also possess these symmetry operations.*

→ Symmetry of the microstructure gives us the form of the effective tensors

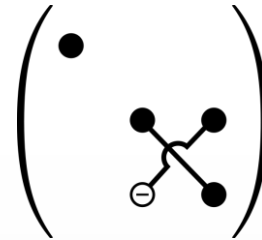
# Symmetry considerations

**Triclinic**



**Tetragonal**

One four-fold rotational axis

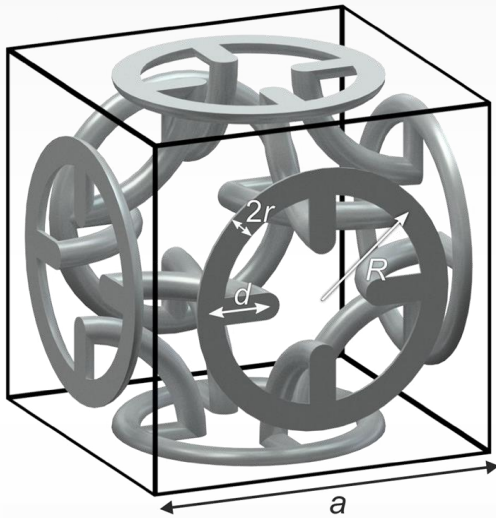


**Cubic**

Four three-fold rotational axes

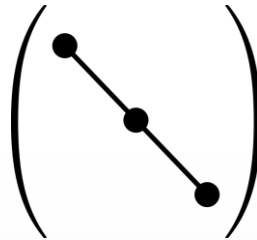


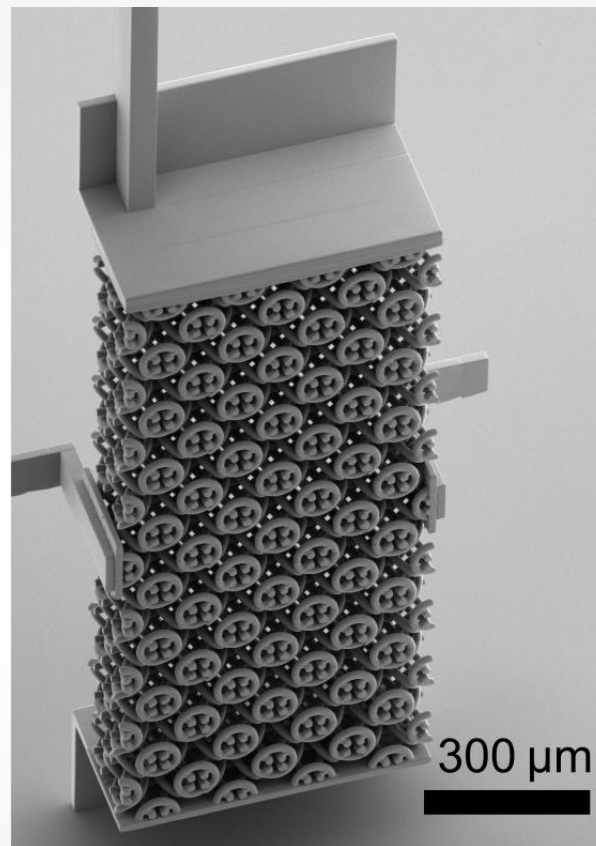
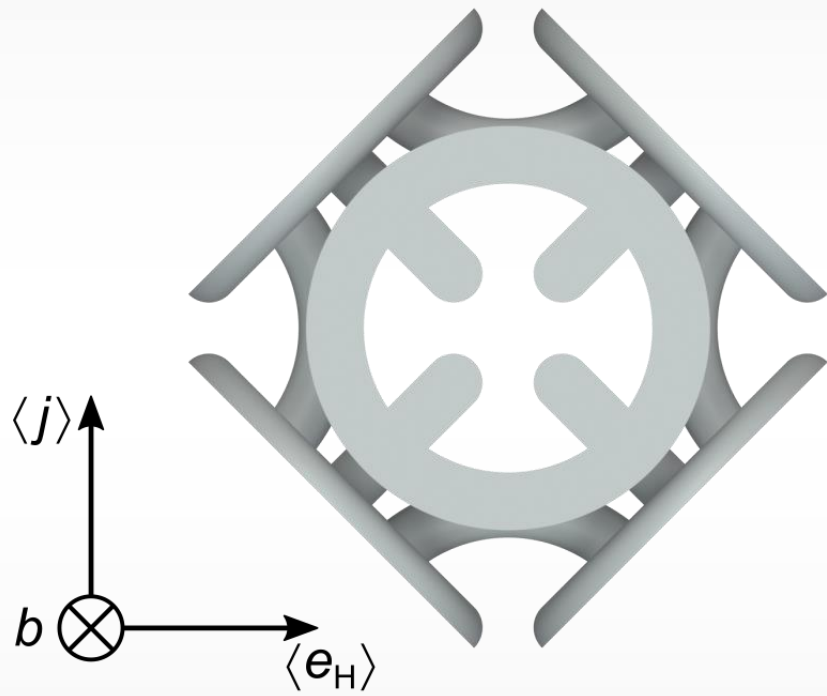
# Symmetry considerations



## Cubic

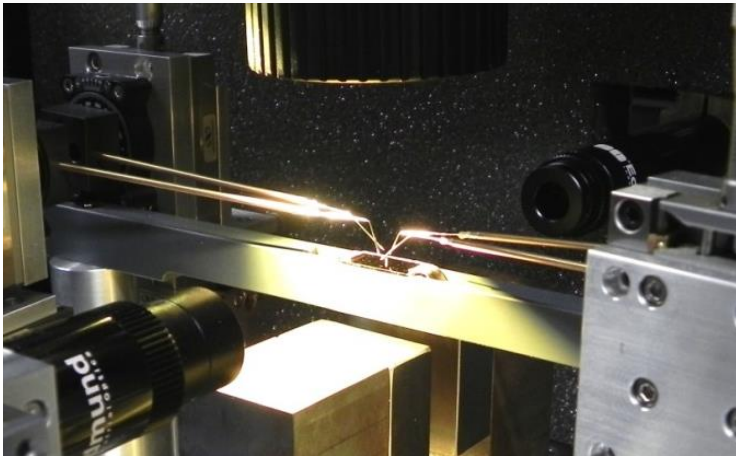
One four-fold rotational axis





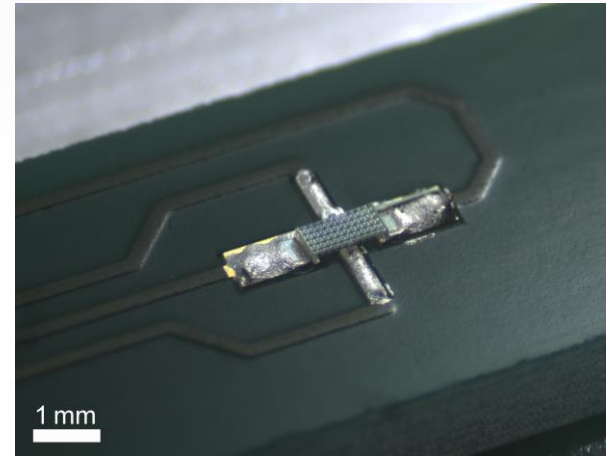
# Current state of experiments

Probe station based



- “Less demanding” fabrication

Integrated device

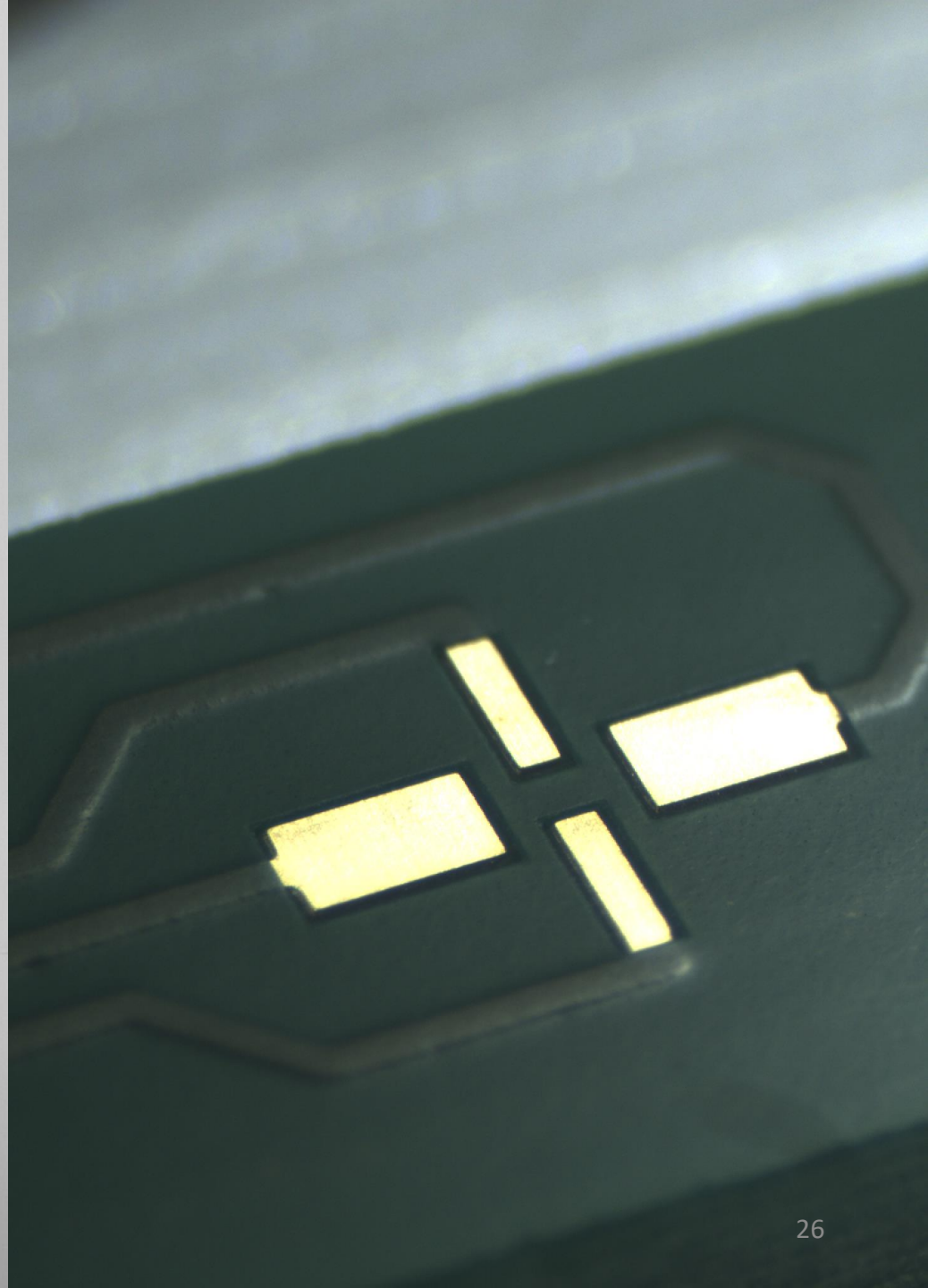
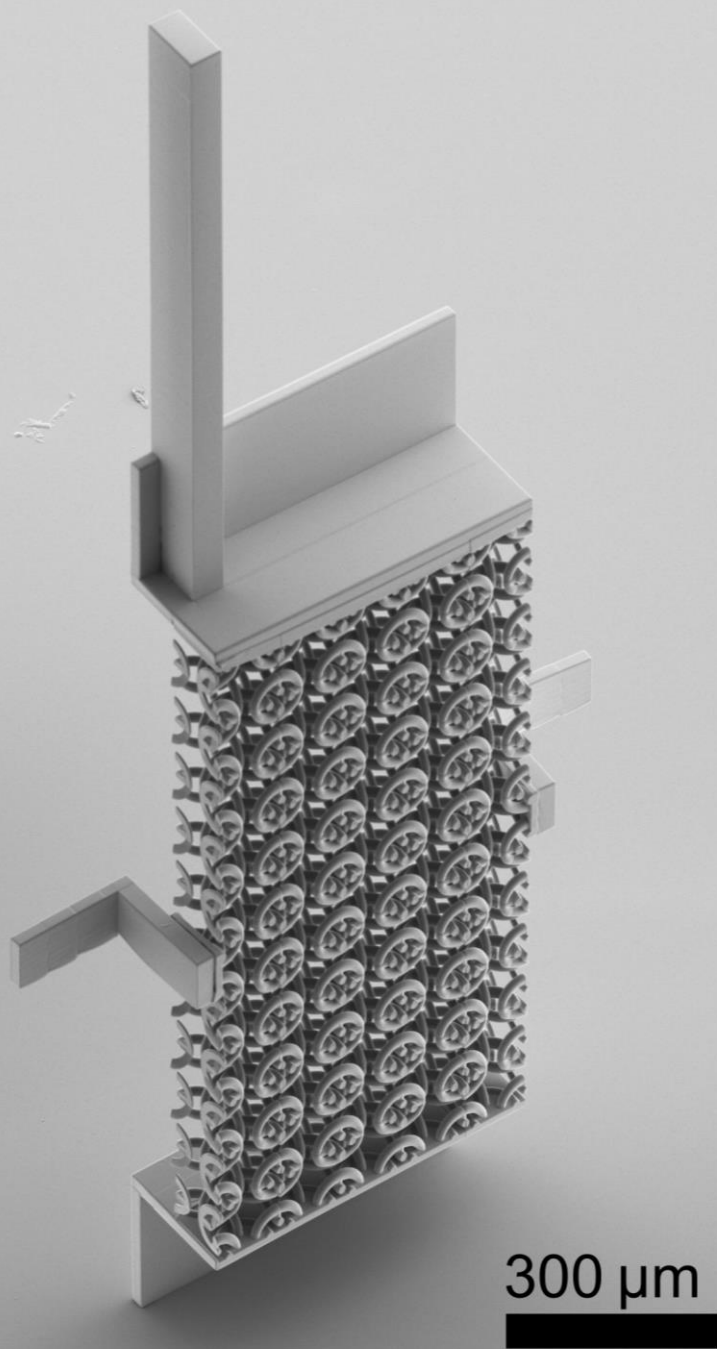


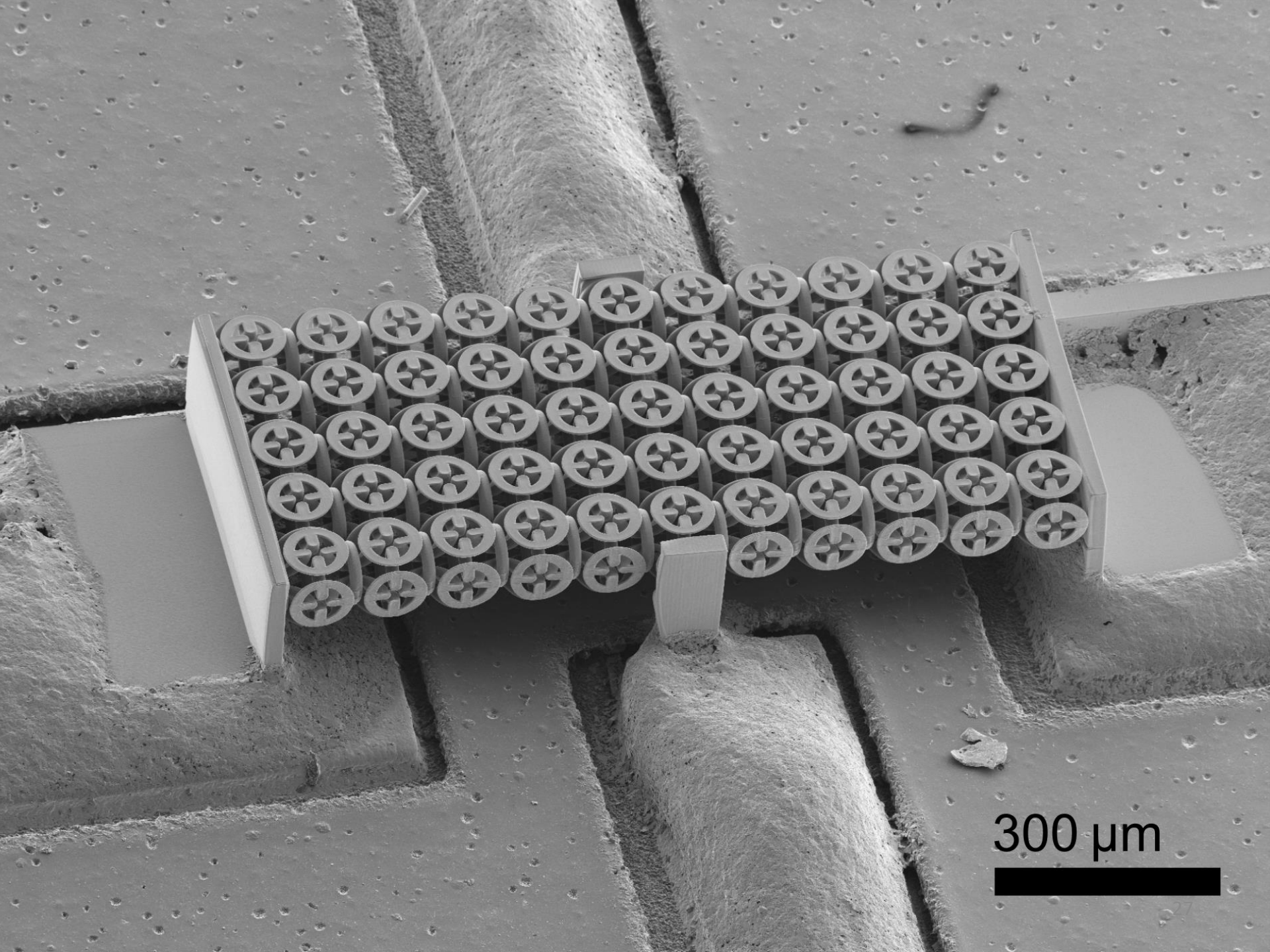
- Facilitates measurements
- Prerequisite for applications



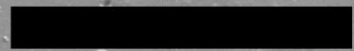


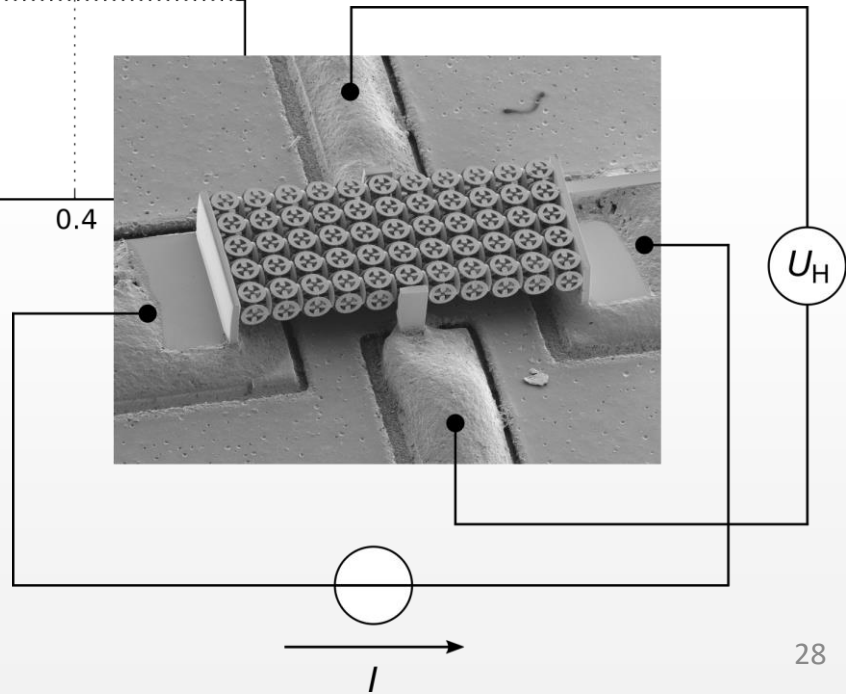
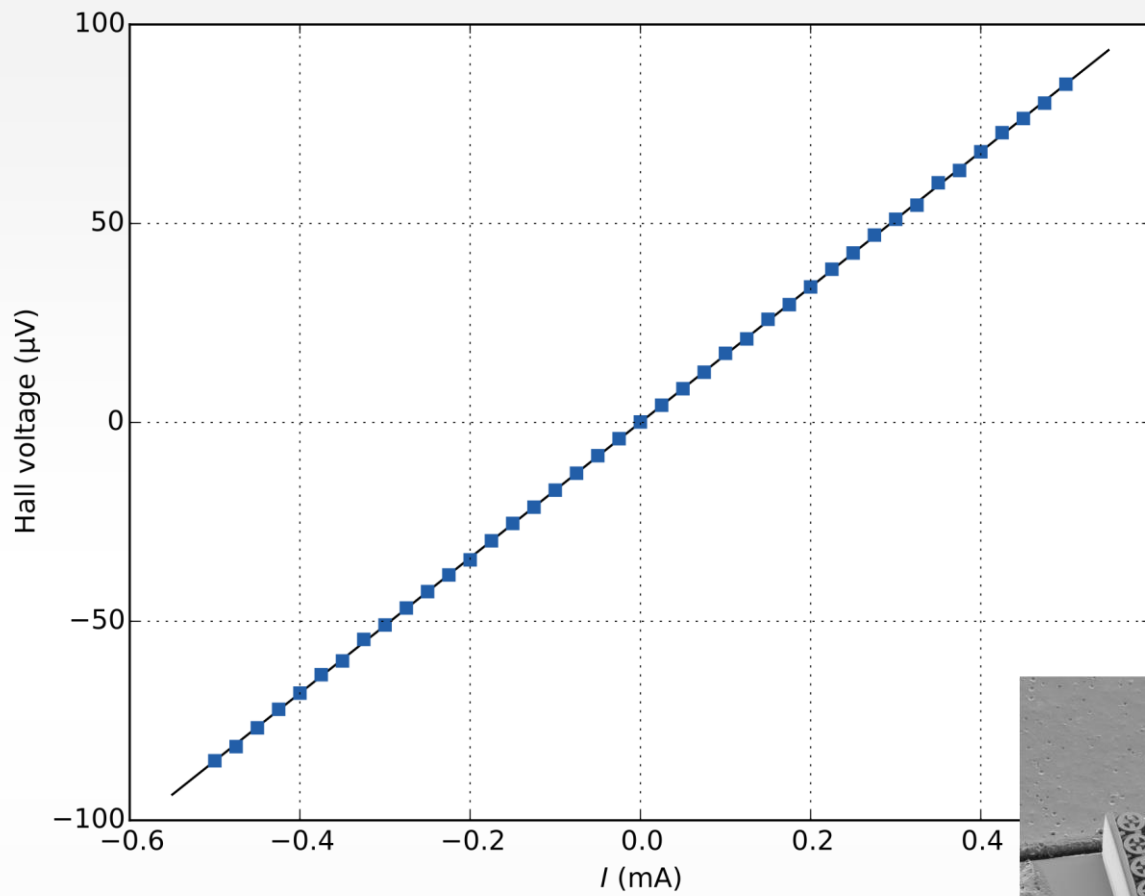
1 mm



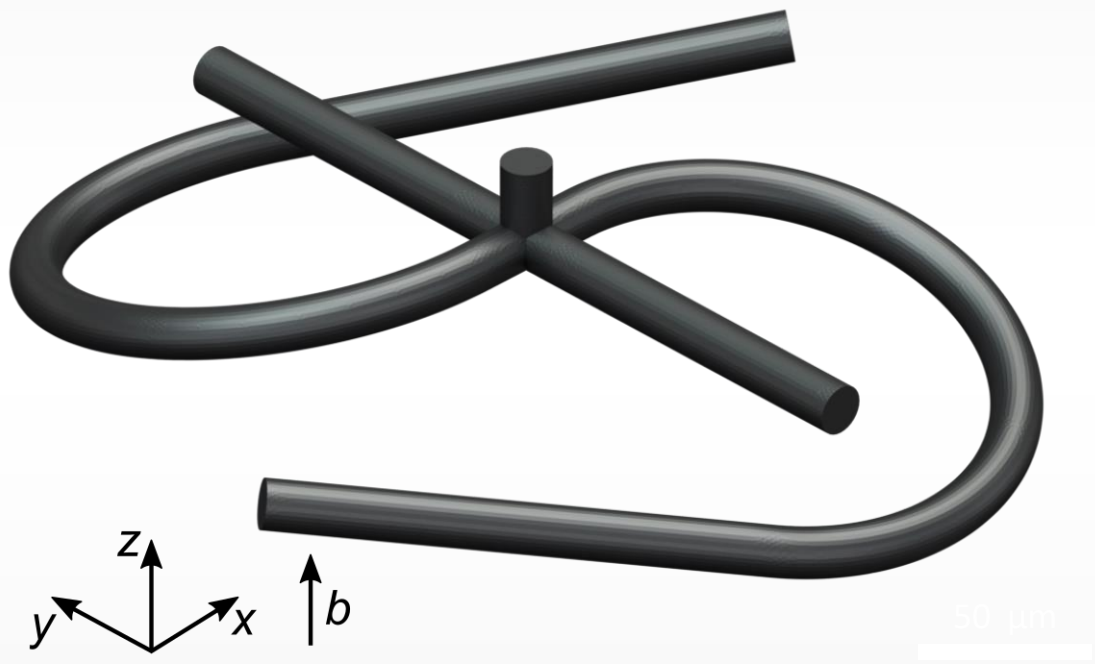


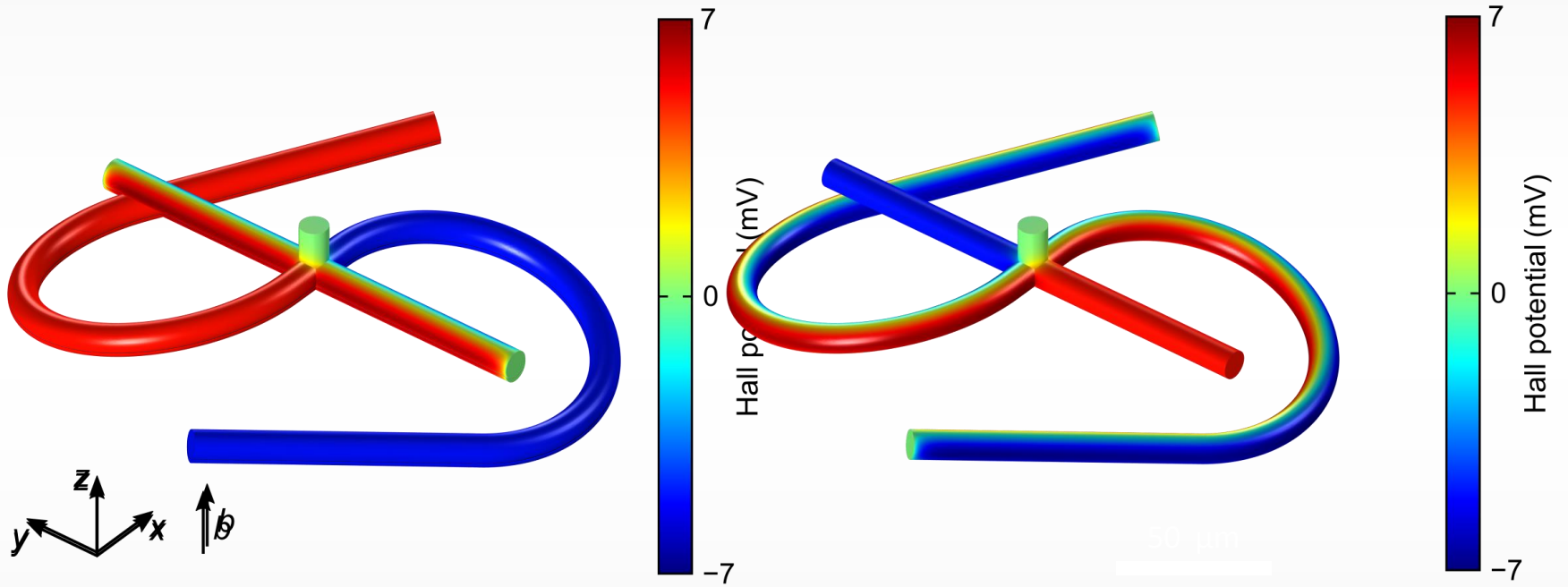
300  $\mu\text{m}$

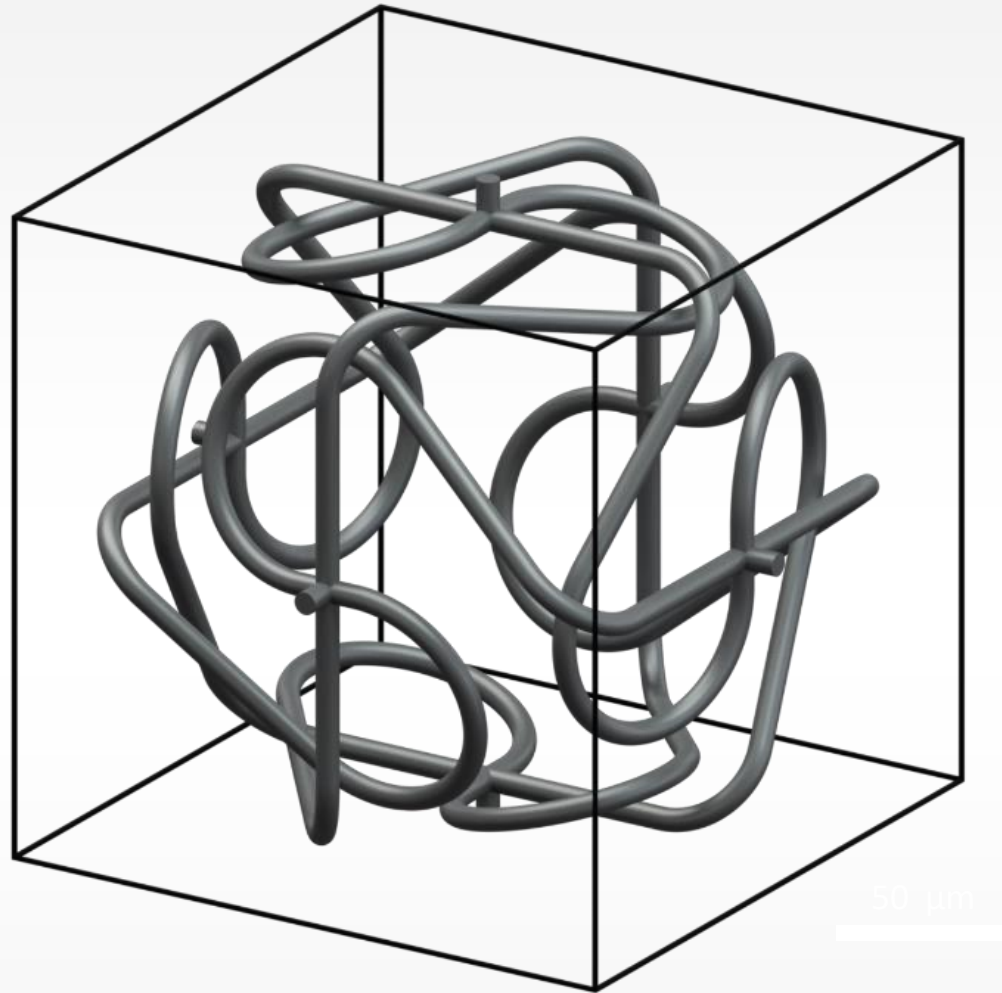




Are there other composites showing a sign-inverted effective Hall coefficient?







$$\text{Cof}(\boldsymbol{\sigma}_0^*) \mathbf{A}_H^* = \langle \text{Cof}(\boldsymbol{\sigma}_0 \nabla \Phi)^\top \mathbf{A}_H \rangle$$



# Anisotropic structures

# An antisymmetric Hall tensor

$$\mathbf{A}_H = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & A_{23} \\ 0 & -A_{23} & 0 \end{pmatrix}$$

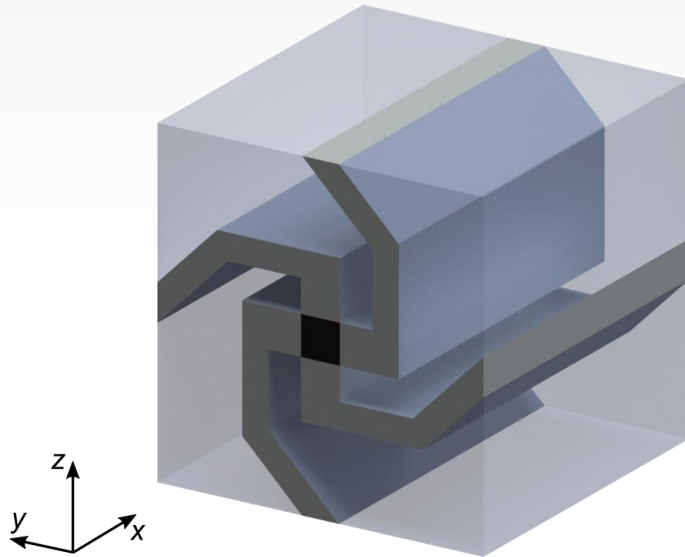
Current along  $\hat{\mathbf{x}}$ , magnetic field in the  $yz$ -plane

$$\mathbf{j} = j_x \hat{\mathbf{x}} \quad \mathbf{b} = b_y \hat{\mathbf{y}} + b_z \hat{\mathbf{z}}$$

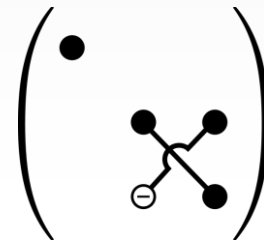
$$\longrightarrow \mathbf{e}_H = \mathcal{E} (\mathbf{A}_H \mathbf{b}) \mathbf{j} = -A_{23} j_x (b_y \hat{\mathbf{y}} + b_z \hat{\mathbf{z}})$$

The Hall electric field is parallel to the magnetic field: ***Parallel Hall effect***

# An antisymmetric Hall tensor

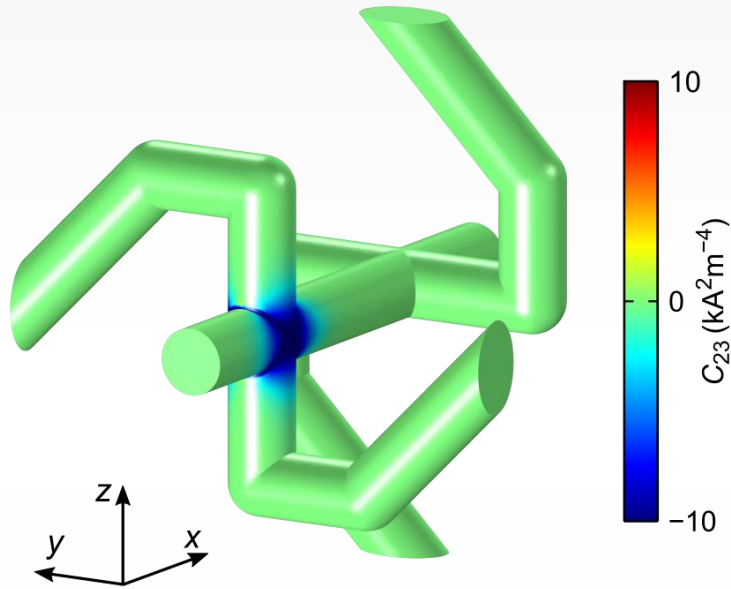


Tetragonal symmetry implies

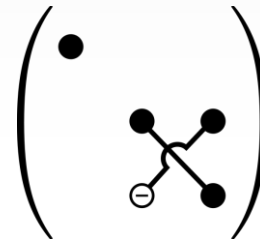


$$\mathbf{A}_{\text{H}}^* = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix} \mathbf{A}_{\text{H}}^0$$

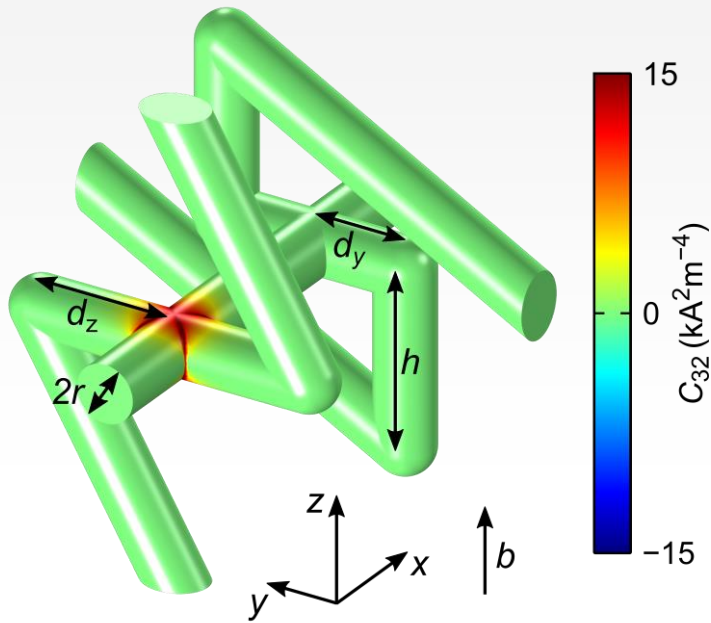
# An antisymmetric Hall tensor



Tetragonal symmetry implies

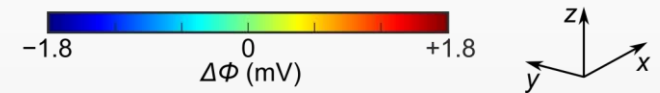
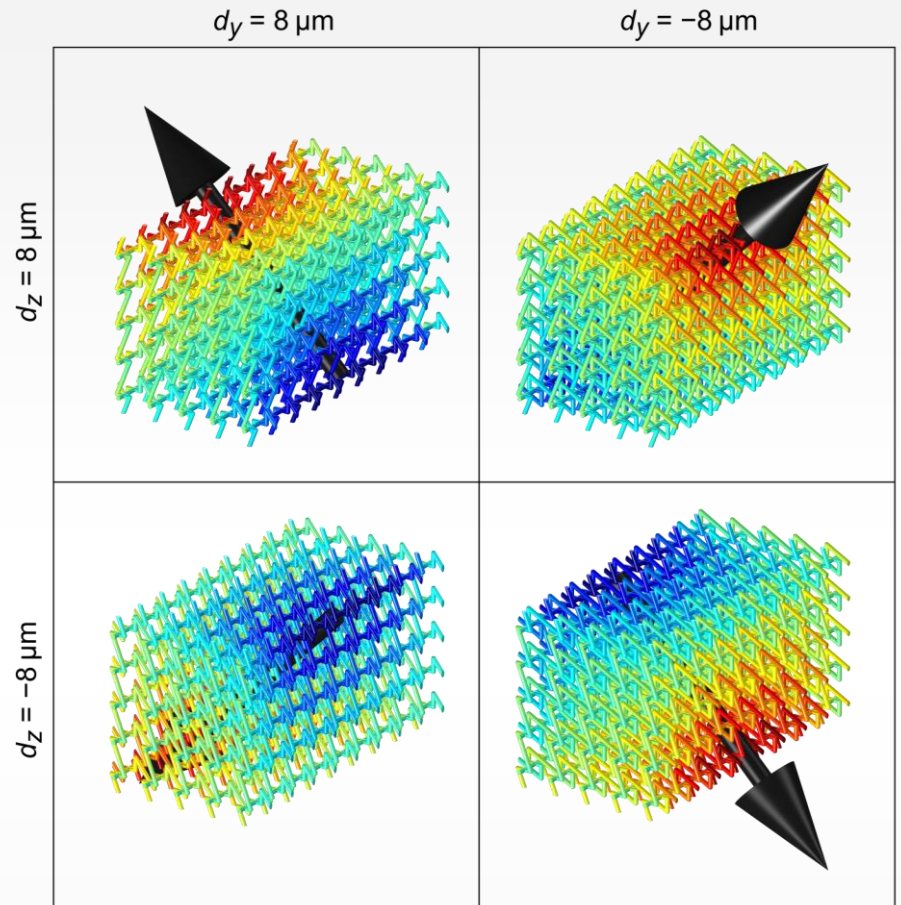


$$\mathbf{A}_H^* = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0.05 & 8.81 \\ 0 & -8.81 & 0.05 \end{pmatrix} \mathbf{A}_H^0$$



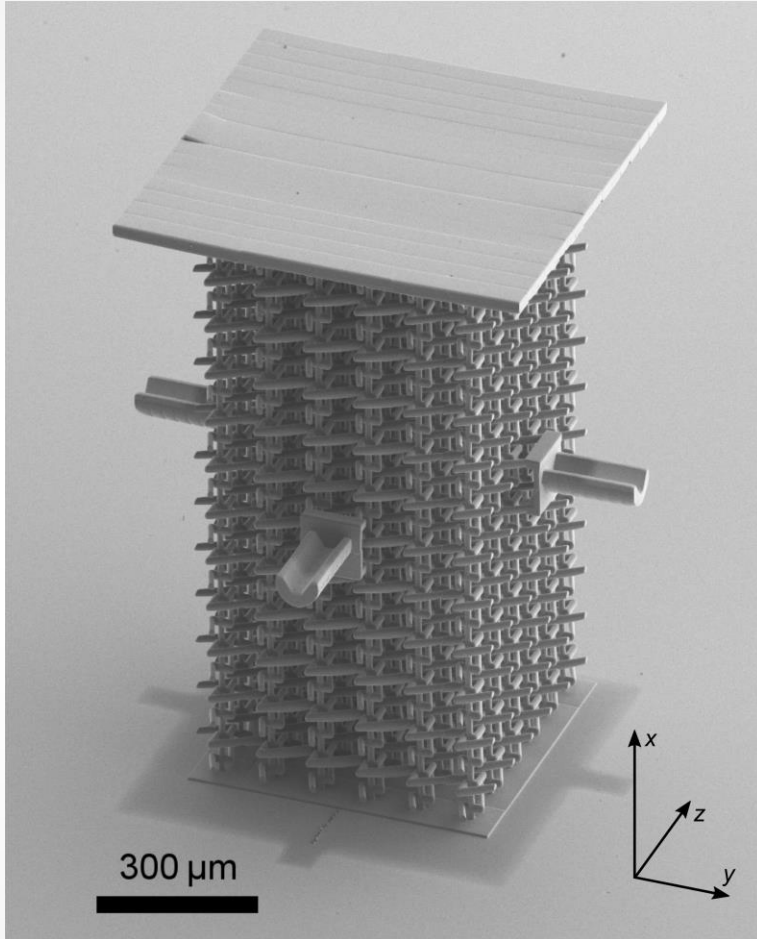
Monoclinic symmetry implies

$$\begin{pmatrix} \bullet & & & \\ & \bullet & & \\ & & \bullet & \\ & & & \bullet \end{pmatrix}$$

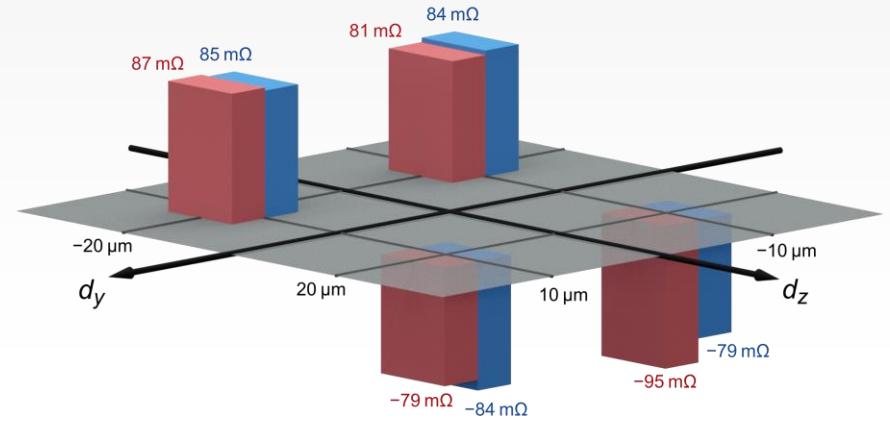


C. Kern *et al.*, Appl. Phys. Lett. **107**, 132103 (2015)

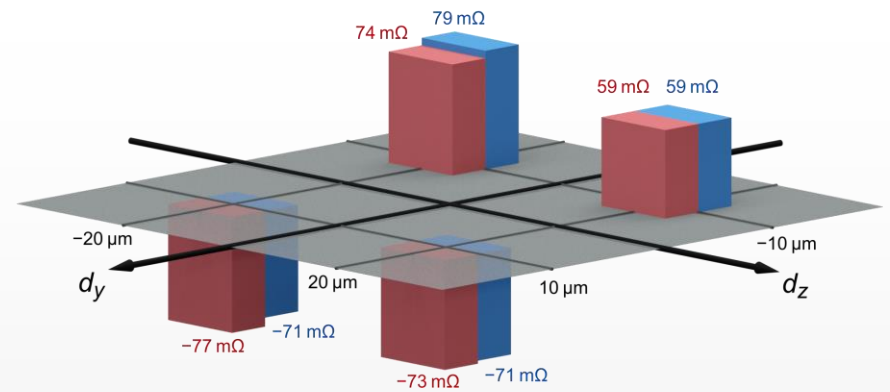
C. Kern *et al.*, arXiv:1806.04914 [cond-mat.mes-hall]

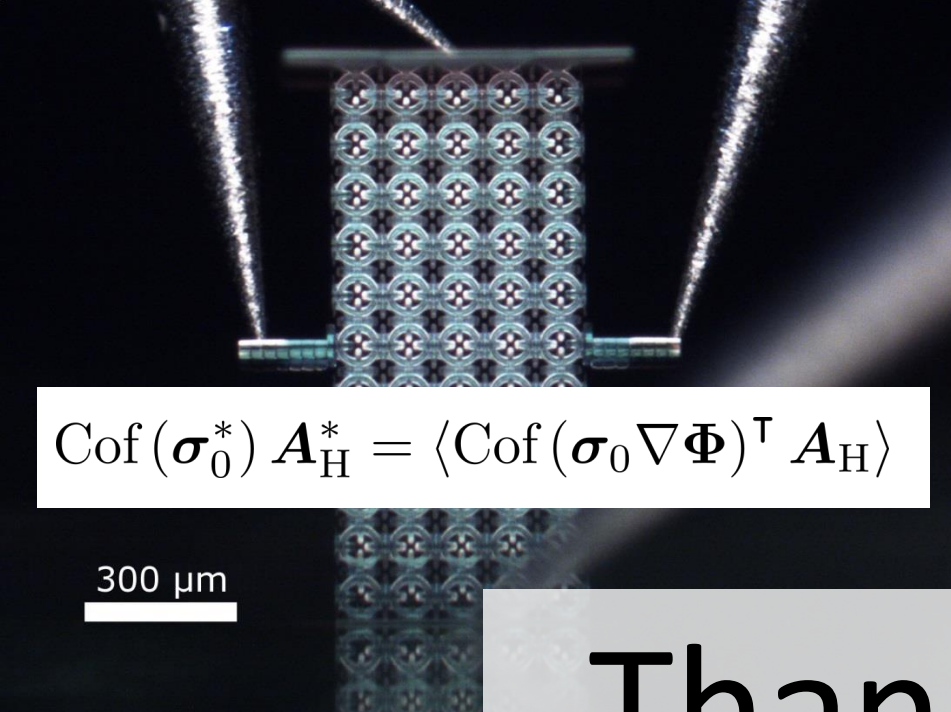


Parallel Hall resistance  $R_H^z$

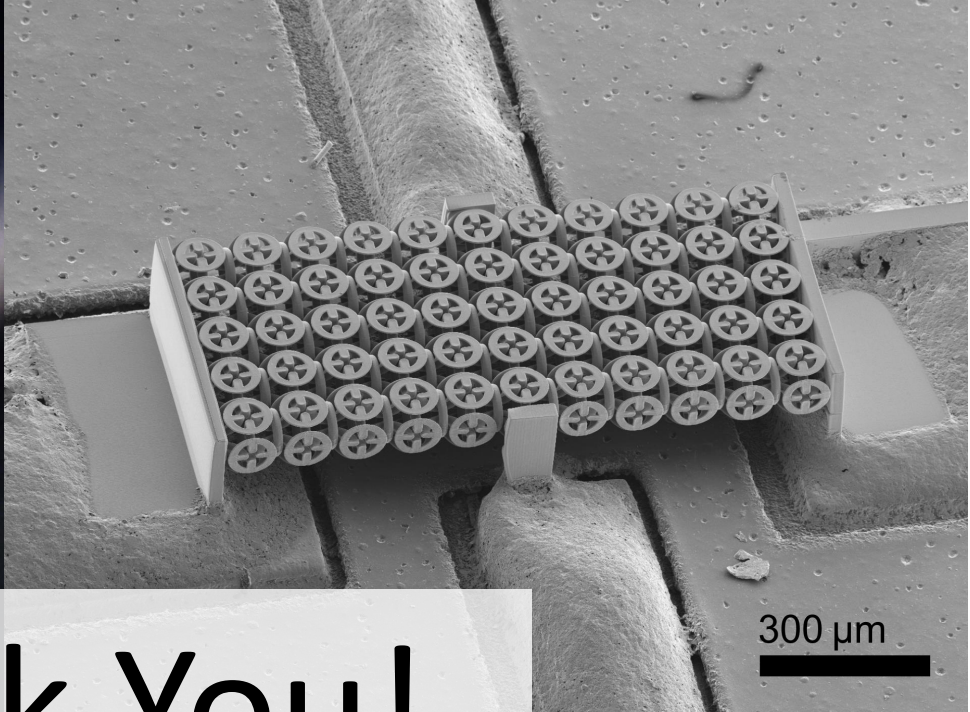


Perpendicular Hall resistance  $R_H^y$





$$\text{Cof}(\boldsymbol{\sigma}_0^*) \mathbf{A}_H^* = \langle \text{Cof}(\boldsymbol{\sigma}_0 \nabla \Phi)^\top \mathbf{A}_H \rangle$$



Thank You!

